

Intra-beam Scattering: benchmarking and predictions

J. Wei, A. Fedotov, W. Fischer, N. Malitsky, G. Parzen et al

BNL

J. Qiang

LBNL

RHIC Beam Experiment Workshop

September 16 - 17, 2004



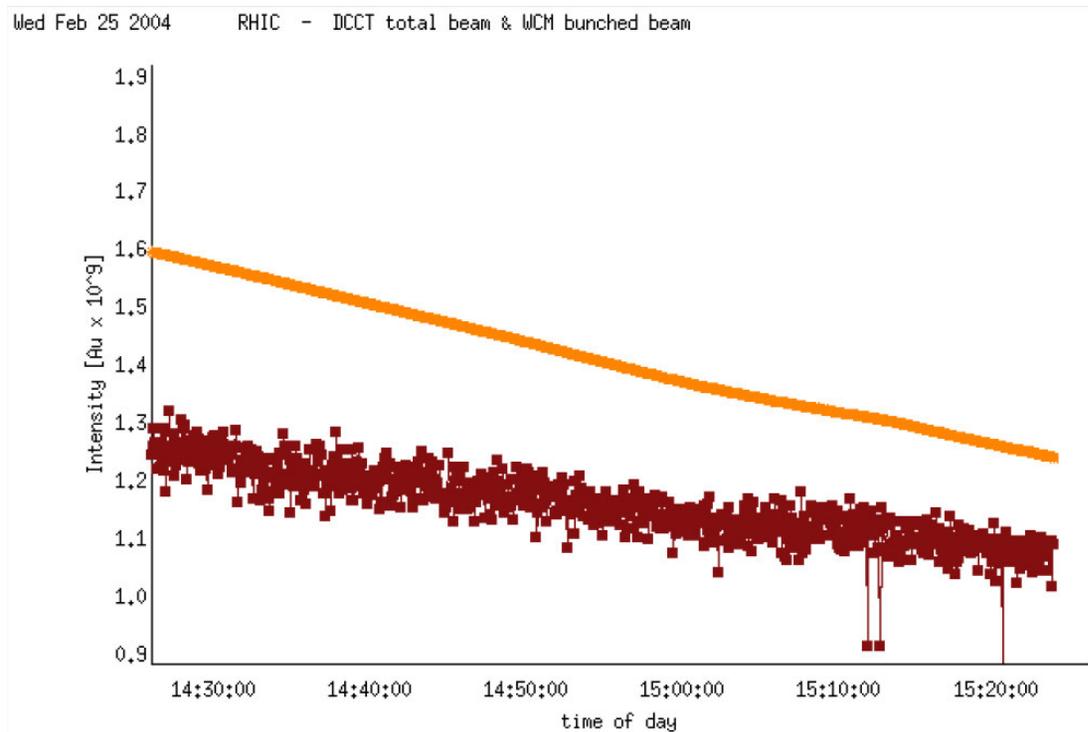
Outline

- Introduction
 - IBS effects, mechanism, scaling laws
 - IBS beam experiment goals
- IBS study results from Run 4
 - beam emittance growth
 - Beam loss & de-bunching
 - Beam distribution evolution: Gaussian-like vs. hollow beams
- IBS predictions for Run 5 Cu operation

- Acknowledgements: M. Blaskiewicz, M. Brennan, R. Connolly, V. Litinenko, T. Satogata, S. Tepikian, J. van Zeijt ...

IBS effects

- Luminosity degradation
 - Transverse emittance growth
 - Longitudinal growth & beam loss due to RF bucket limitation
- De-bunching & experimental background



Mechanism

- Particle motion in the beam rest frame

$$H = \begin{cases} \frac{1}{2} (P_x^2 + P_y^2 + P_z^2) + \frac{1}{2} x^2 - \gamma x P_z + V_C & \text{(bending section)} \\ \frac{1}{2} (P_x^2 + P_y^2 + P_z^2) - \frac{n_1}{2} (x^2 - y^2) + V_C + U_s & \text{(straight section)} \end{cases}$$

- Intra-beam Coulomb scattering among particles of the same bunch

$$V_C = \sum_j \frac{1}{\sqrt{(x_j - x)^2 + (y_j - y)^2 + (z_j - z)^2}}$$

- IBS: Rutherford scattering + Lorentz transformation + averaging

Positive- and negative-mass regime

- Transformed Hamiltonian

$$\bar{H} = \omega_{\beta_x} J_x + \omega_{\beta_y} J_y + \frac{1-\gamma^2 F_z}{2} P_z^2 + V_C \quad \langle F_z \rangle = \frac{1}{\gamma_t^2}$$

$$F_z = \begin{cases} D + DD' + (D')^2 & \text{(bending section)} \\ DD' + (D')^2 & \text{(straight section)} \end{cases}$$

- Below transition: positive-mass regime dominated by “drift”
 - Quasi-equilibrium (equal velocity in beam rest frame)

$$\left\langle \frac{\sigma_x}{\beta_x} \right\rangle \approx \left\langle \frac{\sigma_y}{\beta_y} \right\rangle \approx \frac{\sigma_p}{\gamma}, \quad \gamma \ll \gamma_T$$

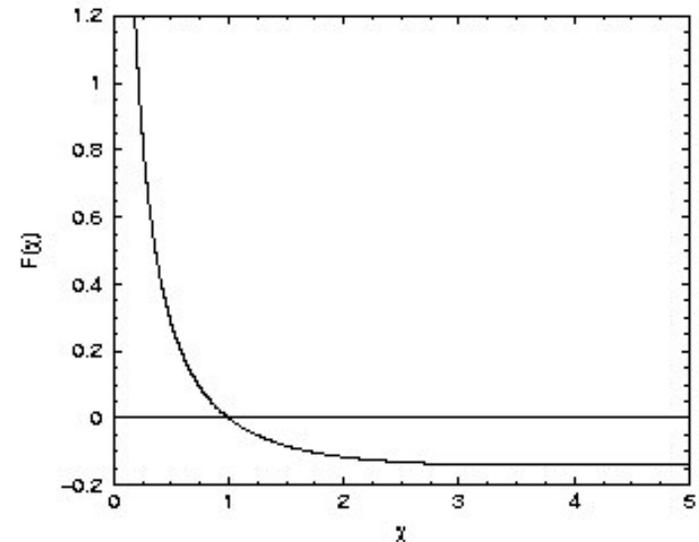
- Temperature exchange between different directions towards equilibrium
 - Actual growth due to AG potential (time-dependent in beam rest frame)
- Above transition: negative-mass regime dominated by diffusion; continuous growth in longitudinal & transverse directions

$$\sqrt{n_b n_c} \langle \sigma_x \rangle \approx \langle D \rangle \sigma_p, \quad \gamma \gg \gamma_T$$

IBS growth scaling

- IBS beam size growth rates assuming an unbounded Gaussian beam

$$\begin{bmatrix} \frac{1}{\sigma_p} \frac{d\sigma_p}{dt} \\ \frac{1}{\sigma_x} \frac{d\sigma_x}{dt} \\ \frac{1}{\sigma_y} \frac{d\sigma_y}{dt} \end{bmatrix} = \frac{Z^4 N}{A^2} \frac{r_0^2 m_0 c^2 L_c}{8\gamma \epsilon_x \epsilon_y S} F(x) \begin{bmatrix} n_b (1 - d^2) \\ -a^2/2 + d^2 \\ -b^2/2 \end{bmatrix}$$



- Proportional to Z^4/A^2
- Proportional to 6-D phase-space density
- Analytic expression for FODO lattice; integral formula for actual lattice
- Inadequate when beam loss occurs / for non-Gaussian beams

Intra-beam scattering study goals

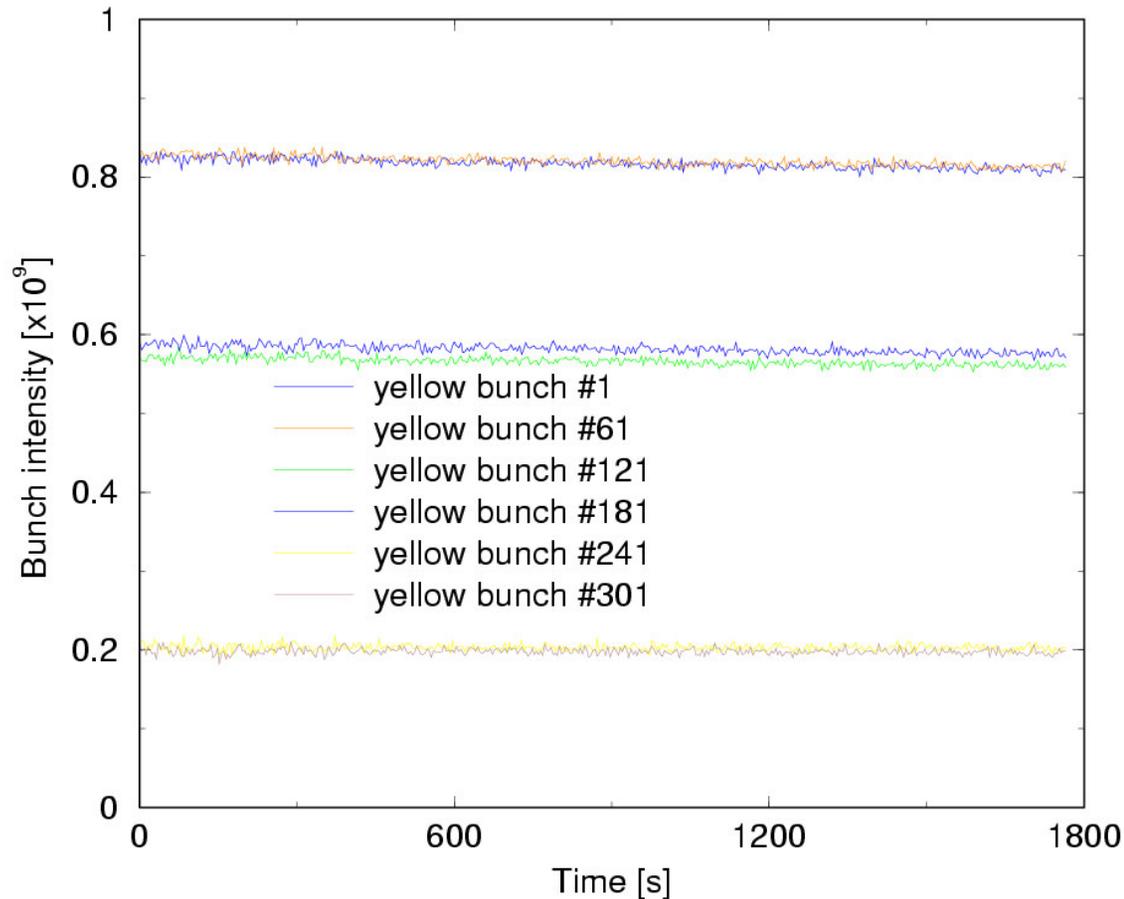
- **Verify the growth of rms beam sizes under IBS**
 - Early theories by Piwinski, Bjorken/Mtingwa
 - Detailed lattice implementation by Martini
 - Asymptotic behavior analysis/ approximation by Parzen
 - Approximation model by Wei
 - Recent compilation by the Russian collaborators / Fedotov
 - **Verify the beam de-bunching behavior under IBS**
 - Predictions by Wei using the Fokker-Planck approach
 - **Verify the longitudinal bunch profile evolution under IBS**
 - Predictions by Wei using the Fokker-Planck approach
 - Similar approach used in stochastic cooling analysis/predictions
- **Help to determine RHIC II performances**

IBS observables

- rms beam sizes (bunch length, transverse emittances)
 - Many early studies performed
 - Need to single out IBS from other processes
 - (beam-beam, tune kicker, Landau cavity, dual RF, RF noise ...)
 - Need to calibrate Ionization Profile Monitor readings
- Beam loss
 - Wall Current Monitor and DCCT readings
 - De-bunching with a single vs. a dual RF (28 MHz and 200 MHz)
- Beam profiles (longitudinal)
 - Hollow bunch profile evolution
 - Gaussian-like bunch profile evolution & asymptotic shape

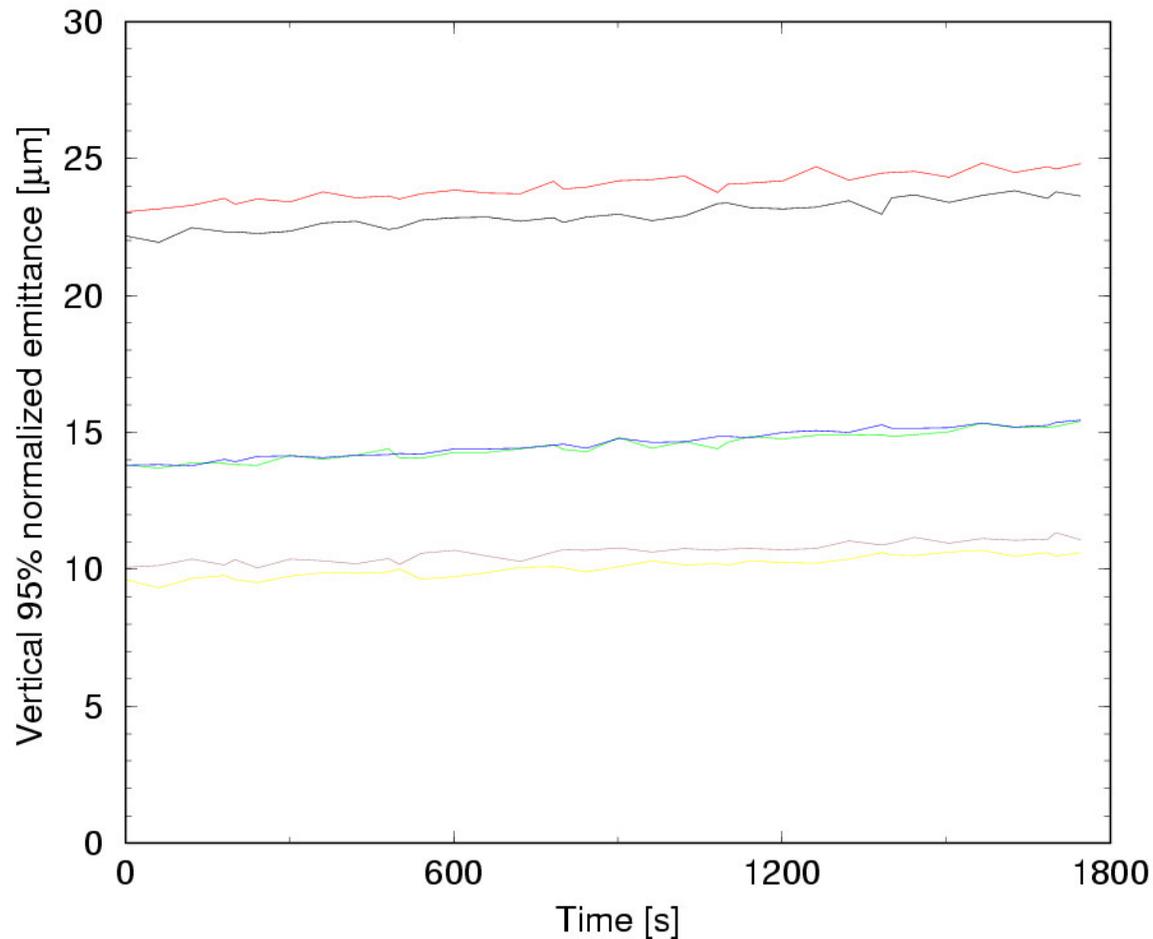
Dedicated IBS studies during Run 4

- January - March, 2004
- Simultaneous IBS measurement under different intensities



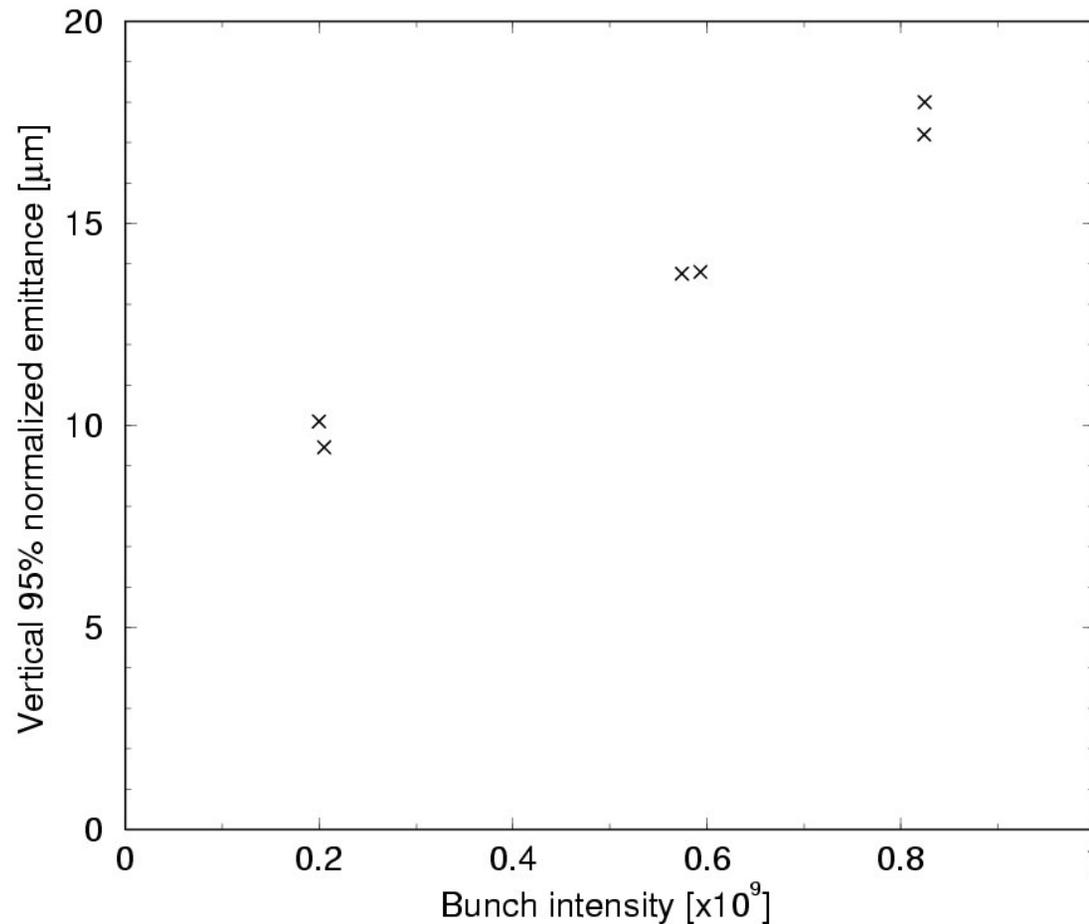
Observation of emittance growths

- Initial transverse emittance depends on intensity
- Overall scaling agrees with IBS theory (Fedotov)



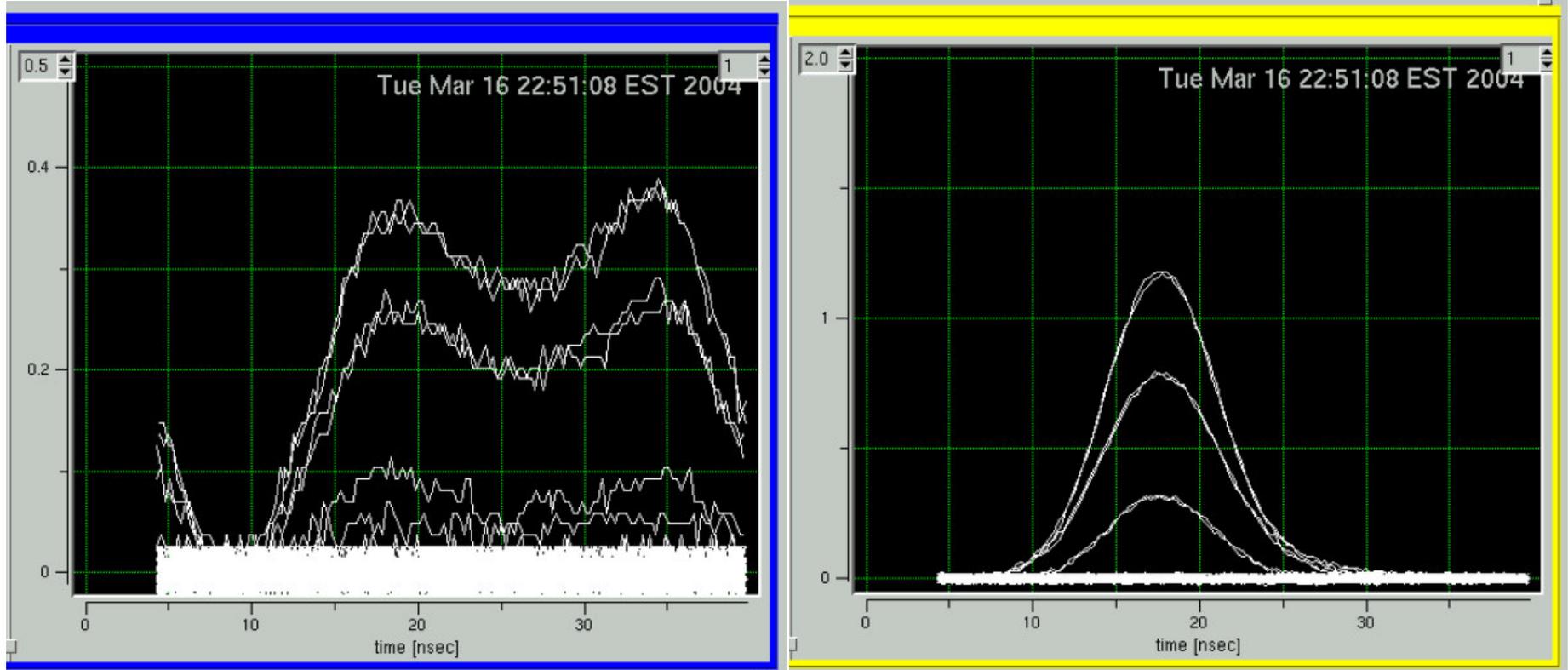
Intensity dependence of emittances

- Possible space-charge effects at Booster/AGS
- Measurement in AGS can confirm IPM calibration in RHIC



March 16, 2004: run #4790

- Gold beam, store at 100 GeV/u with h=360 RF system; no beam collision
- No Landau cavity, no dampers, no kickers
- Hollow beam in blue (phase jump), normal beam in yellow

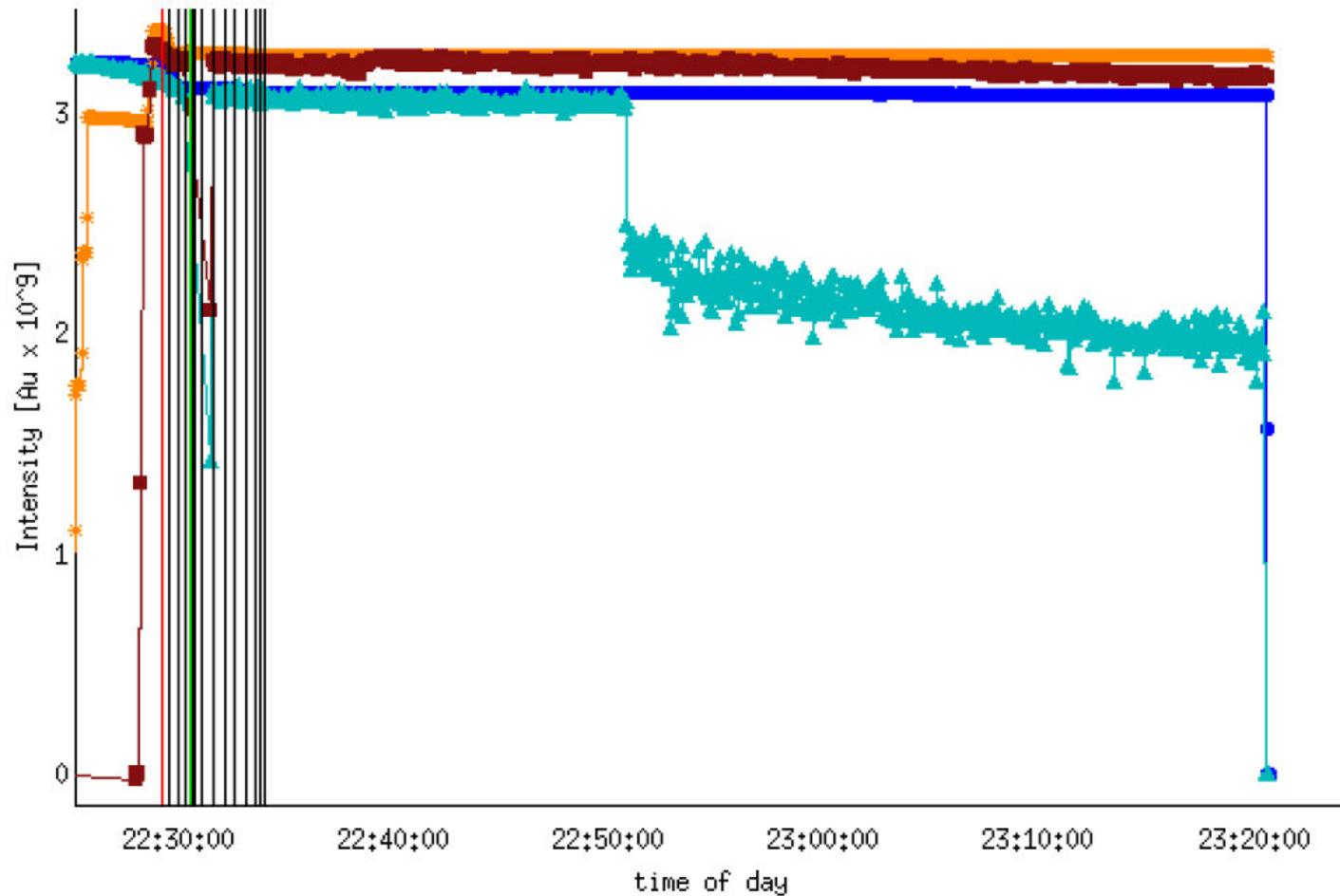


September 16-17, 2004

Wall current monitor intensity readings

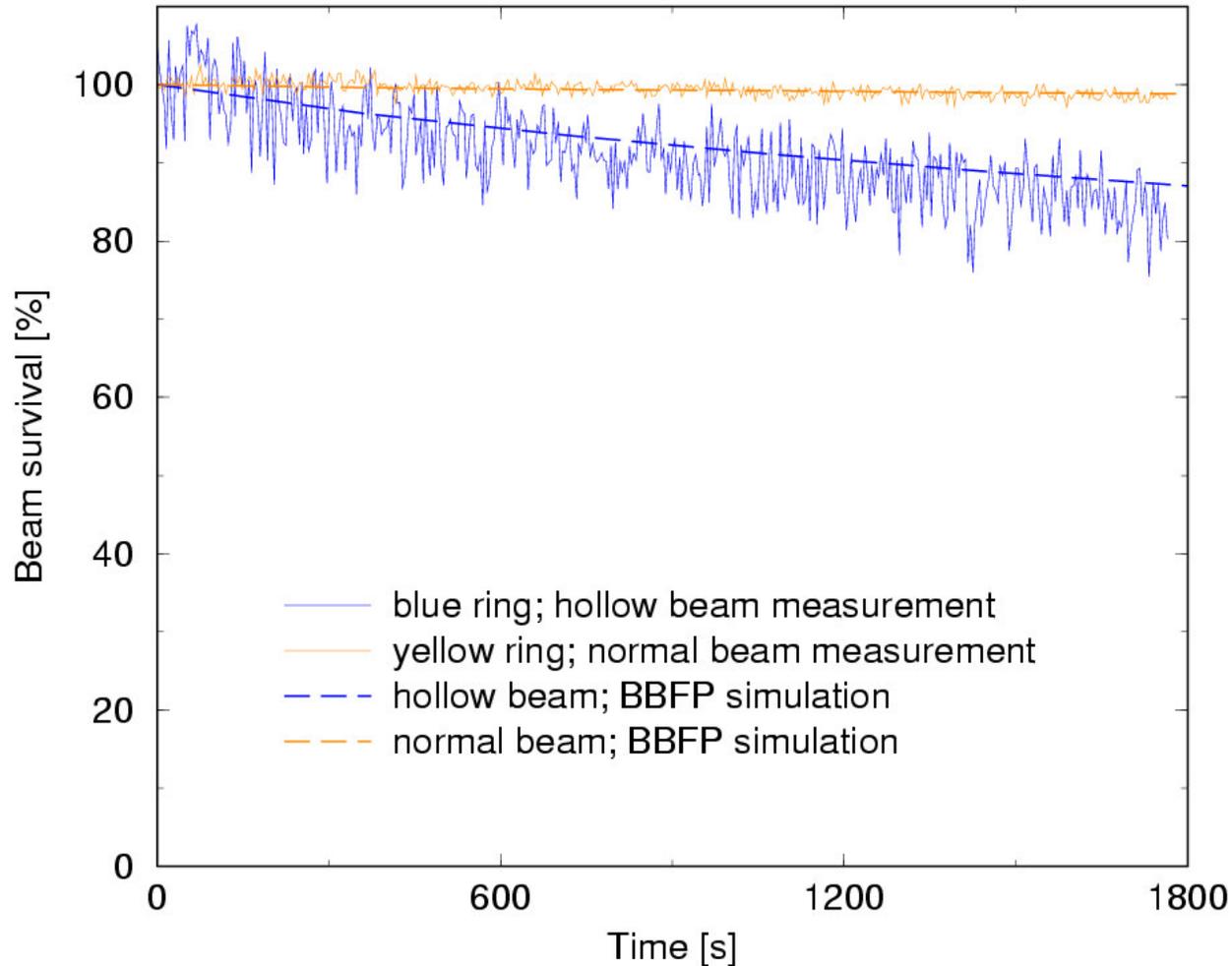
- Distinctively different beam loss (de-bunching) behavior

Tue Mar 16 2004 RHIC - DCCT total beam & WCM bunched beam



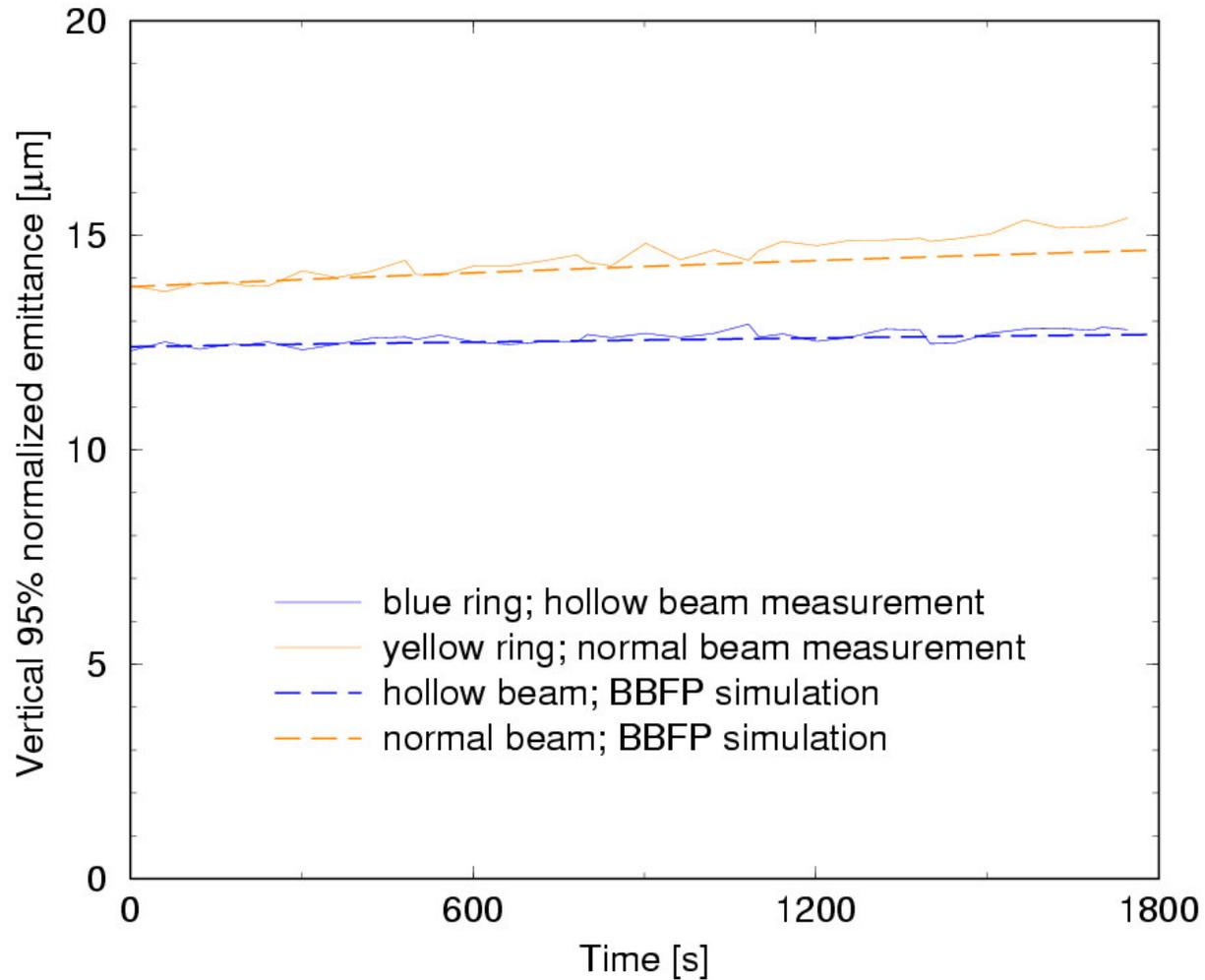
Beam de-bunching loss benchmarking

- Run #4790, 30 minute WCM measurement vs. BBFP code simulation



Transverse emittance growth benchmarking

- Run #4790, 30 minute IPM measurement vs. BBFP code simulation

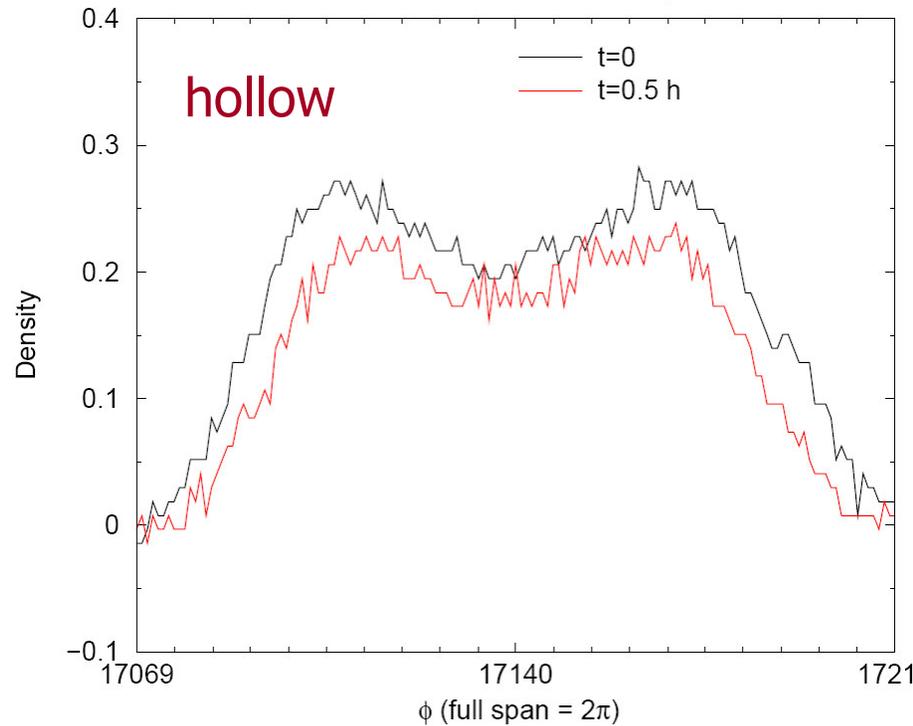


Processed WCM data

- Normal beam: Gaussian-like shape
- Hollow beam: reducing depth of the hole \rightarrow approaching Gaussian

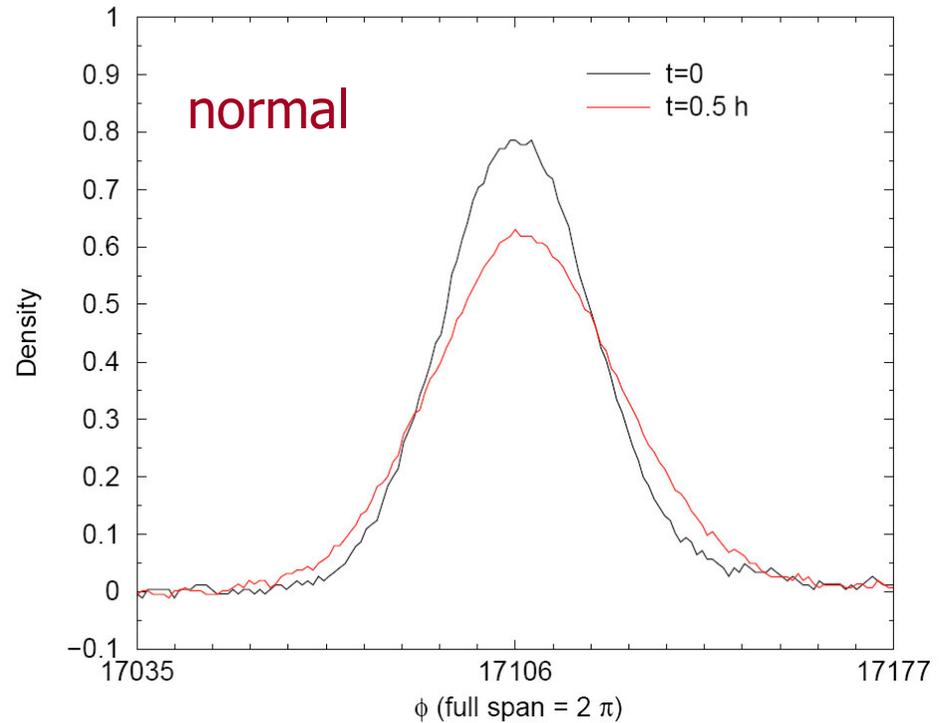
RHIC WCM 2004-3-16 run #4790 Yellow

trace 1079495460 and 1079497224, bunch #3



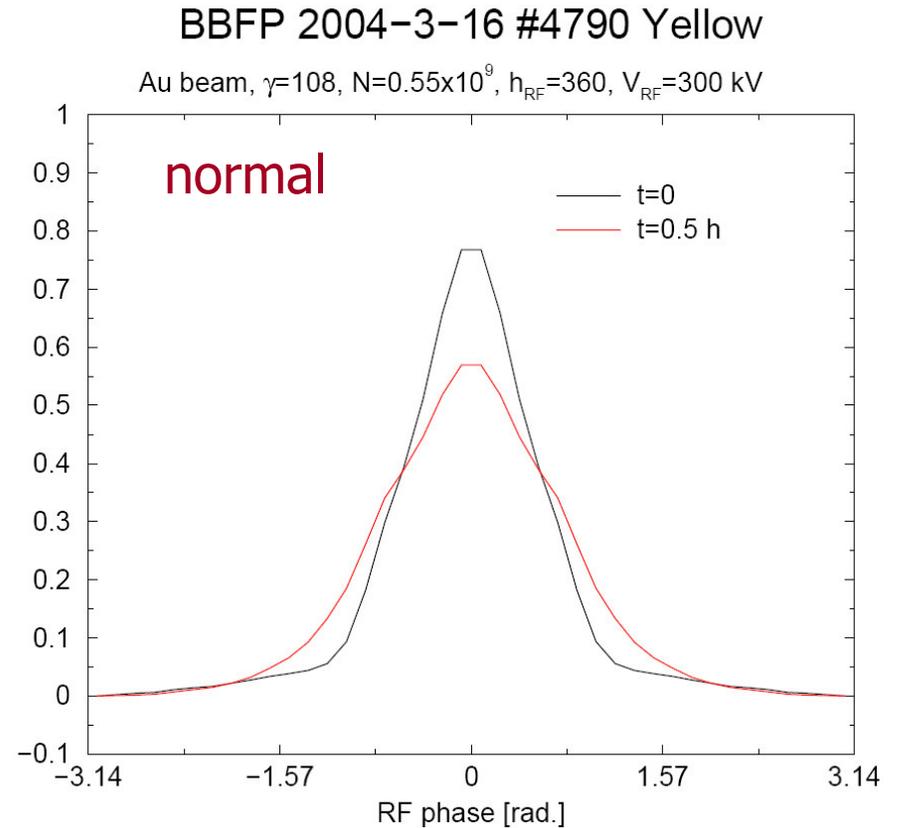
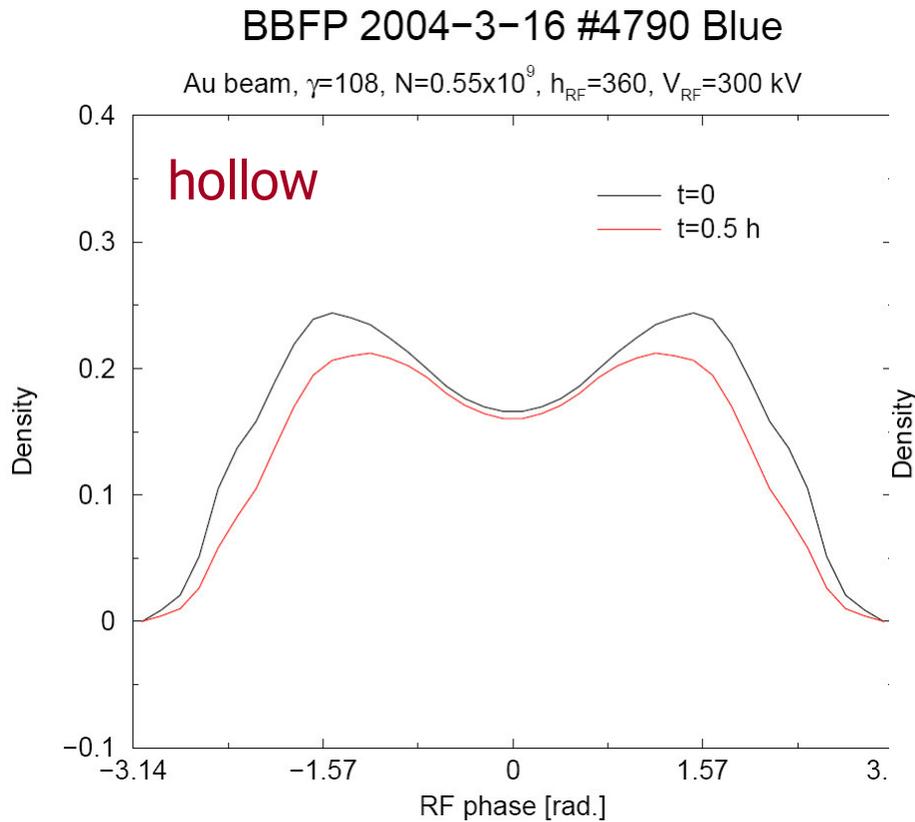
RHIC WCM 2004-3-16 run #4790 Blue

trace 1079495460 and 1079497224, bunch #3



BBFP simulation of the beam profiles

- Good agreement
- Details to be refined



BBFP: Fokker-Planck equation solver

- Drift and diffusion

$$\frac{\partial \Psi}{\partial t} = -\frac{\partial}{\partial J} (F\Psi) + \frac{1}{2} \frac{\partial}{\partial J} \left(D \frac{\partial \Psi}{\partial J} \right), \quad \text{with} \quad \begin{cases} J = 0: & -F\Psi + \frac{D}{2} \frac{\partial \Psi}{\partial J} = 0, \\ J = J_{\text{max}}: & \Psi = 0, \end{cases}$$

where the drift coefficient is given by

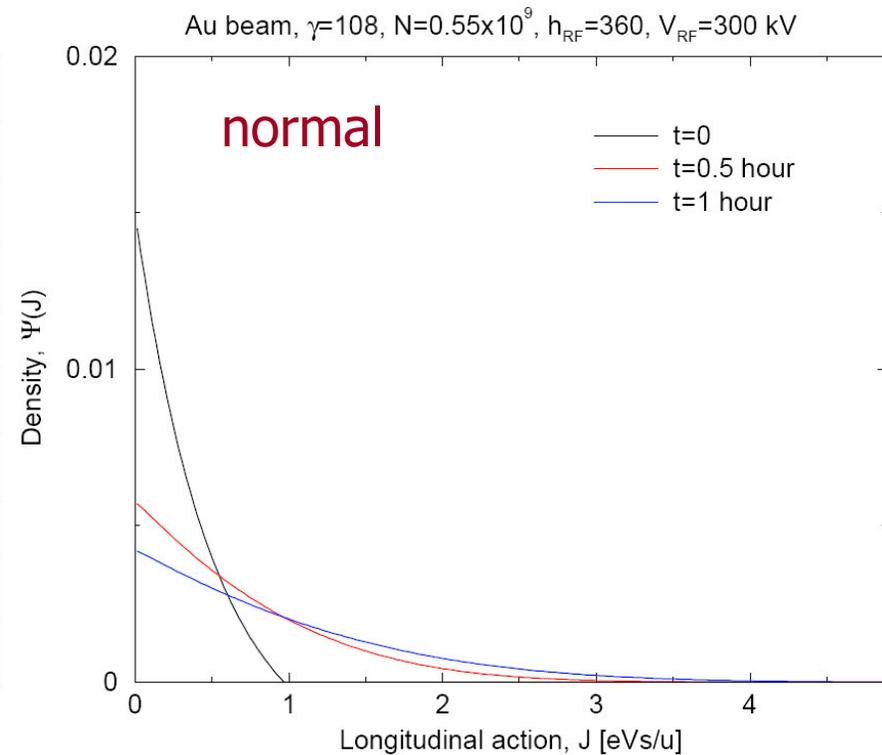
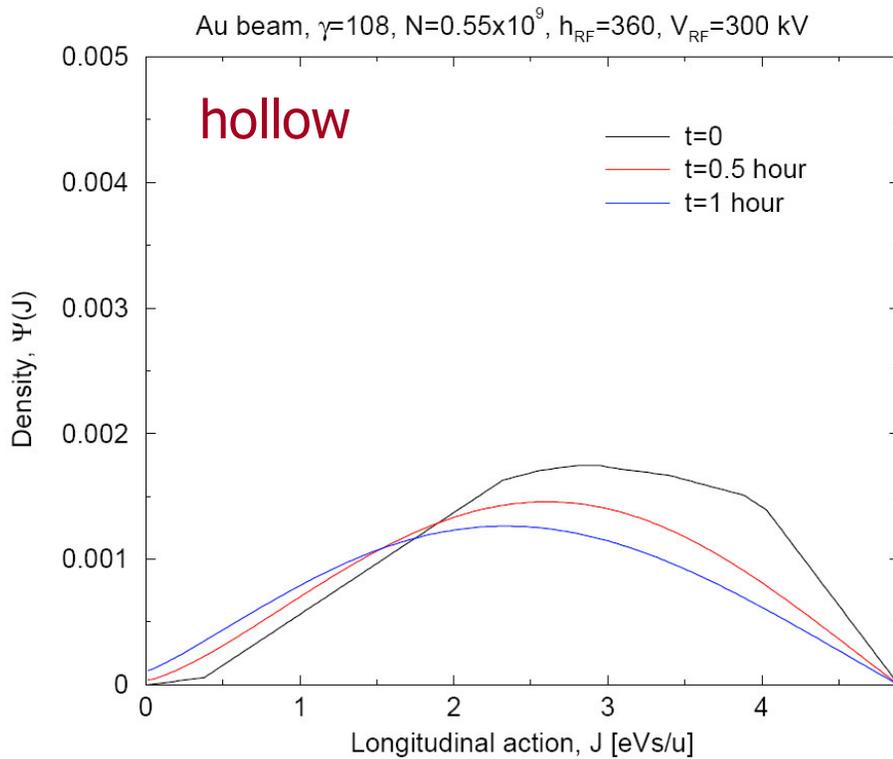
$$F(J) = \oint \frac{2dz}{\pi R} \int_0^{\frac{1}{2}} dQ \left. \frac{\partial W}{\partial J} \right|_{\phi}^{-1} (Q, J) \int_{J_{\text{min}}}^J \left. \frac{\partial W}{\partial J} \right|_{\phi} (Q', J') [A_F(\lambda_1) + A_F(\lambda_2)] \Psi(J') dJ'$$

and the diffusion coefficient is given by

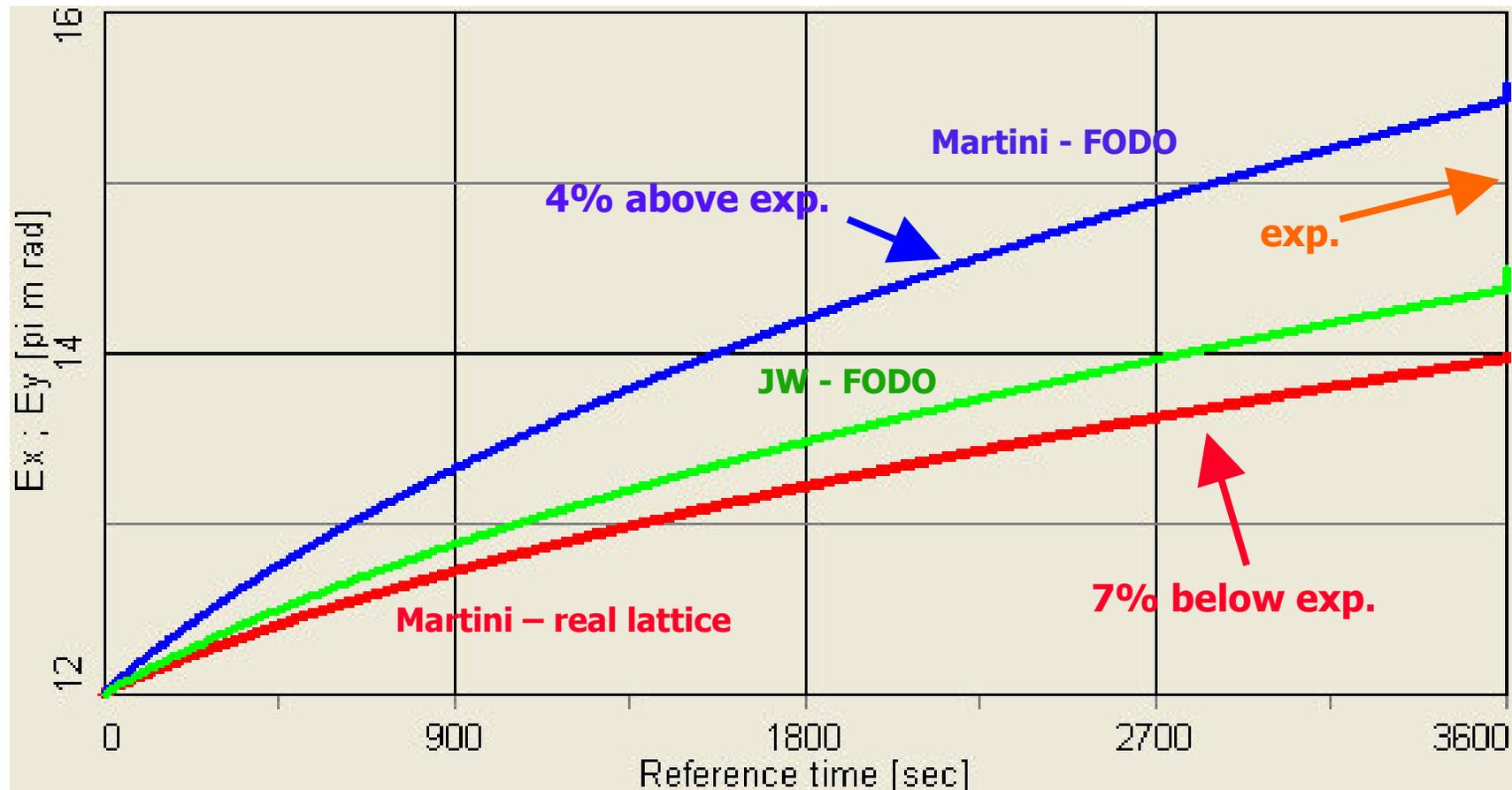
$$D(J) = \oint \frac{2dz}{\pi R} \int_0^{\frac{1}{2}} dQ \left[\left. \frac{\partial W}{\partial J} \right|_{\phi}^{-1} (Q, J) \right]^2 \int_{J_{\text{min}}}^J \left. \frac{\partial W}{\partial J} \right|_{\phi} (Q', J') [A_D(\lambda_1) + A_D(\lambda_2)] \Psi(J') dJ'$$

BBFP simulation of beam evolution

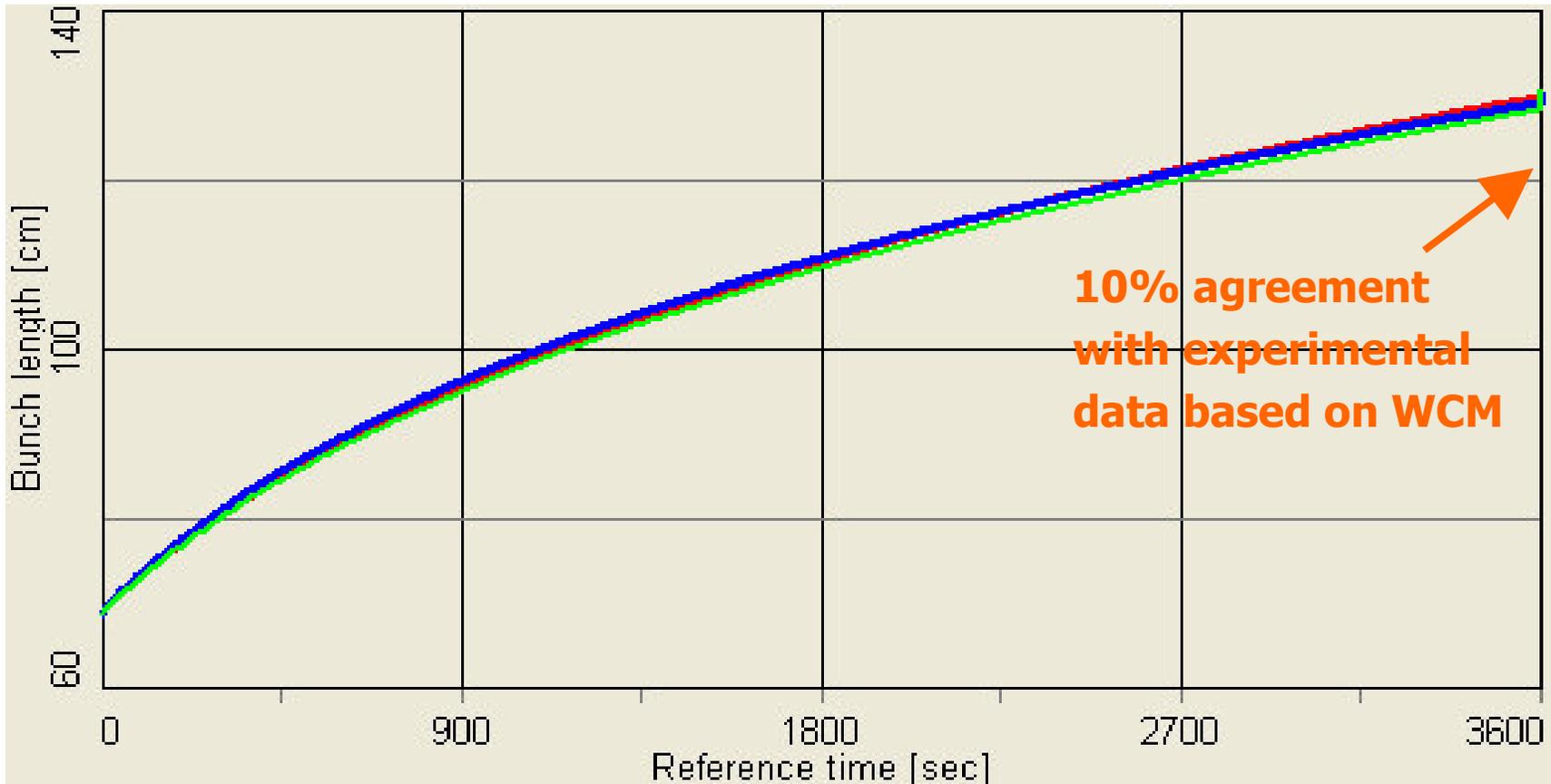
- Density projection in longitudinal action
- Convertible to the phase / momentum planes



Transverse emittance & lattice dependence



Run #4790 - rms bunch length growth



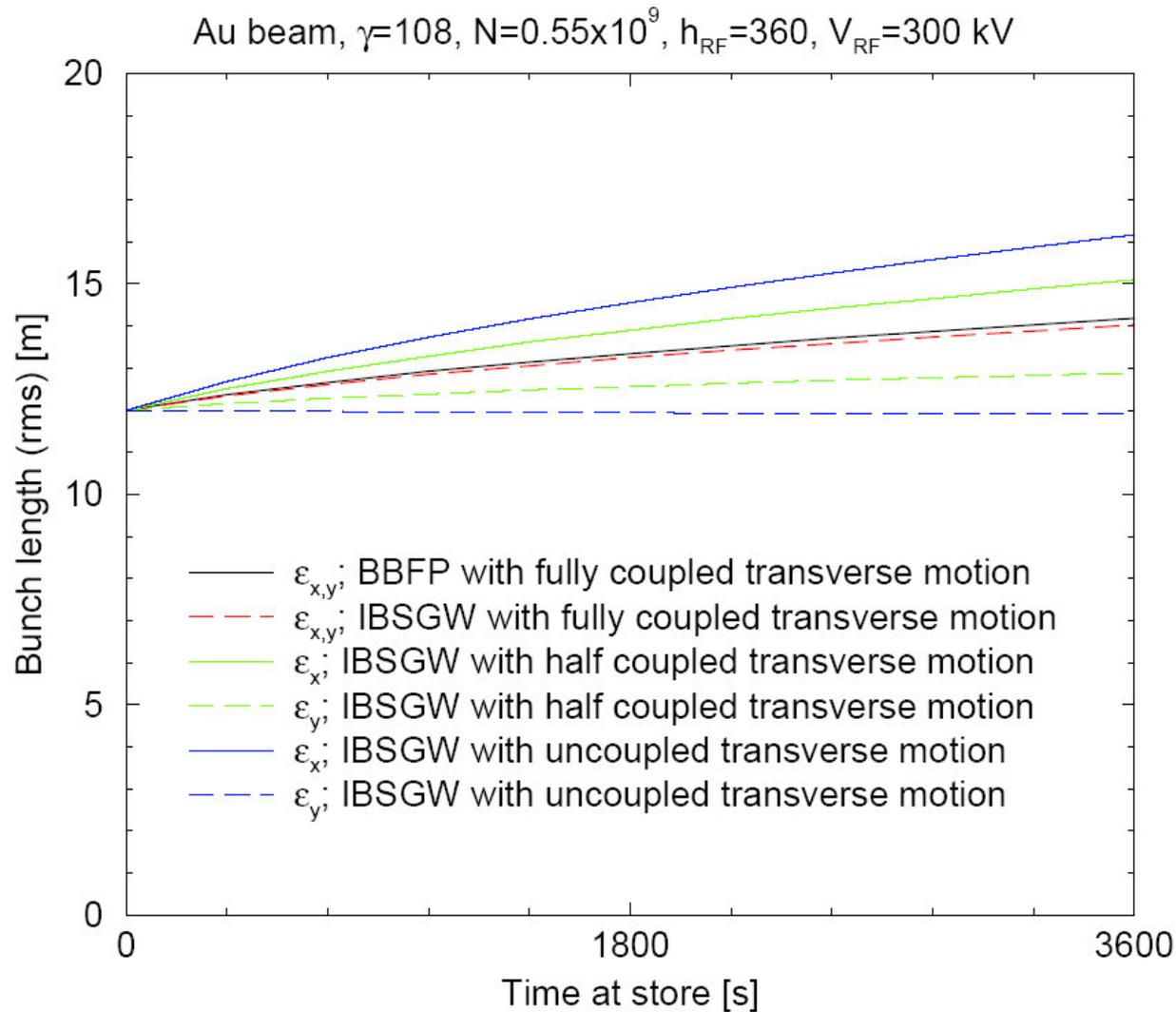
Transverse emittance comparison issues

- Strong dependence on average dispersion
 - Need an accurate estimate of dispersion and dispersion wave
 - Asymptotically $\sigma_x^2 \approx D_p^2 \sigma_p^2$
(Johannes promised to get an estimate from online model)
- Calibration of IPM at store
 - Calibration was done only for lattice at injection
 - Possible IPM electronics degradation/peak suppression may result in falsely large measured emittance value
- Verify source of intensity dependence in Booster / AGS

Puzzles and clues

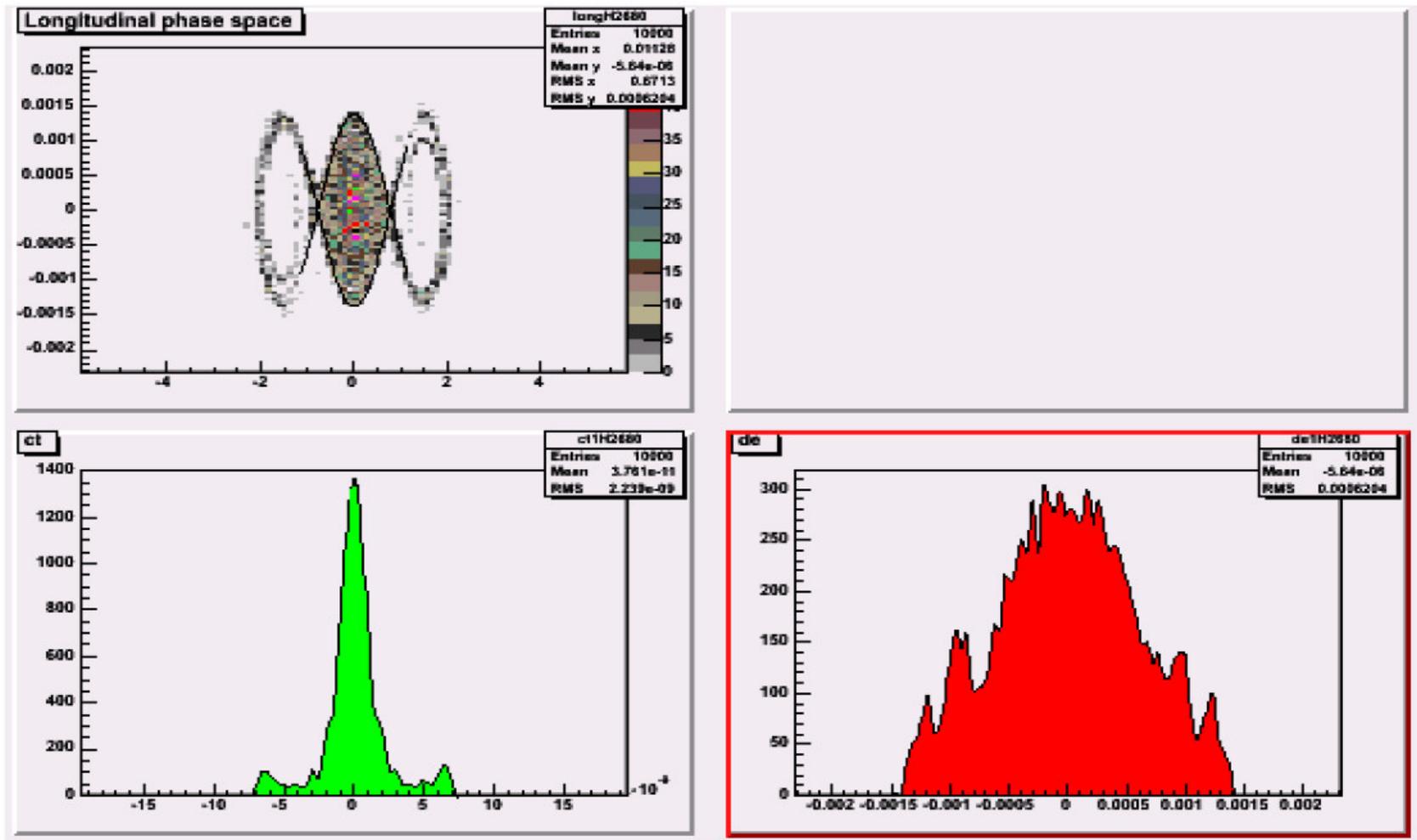
- **Transverse emittance under-estimation**
 - Simulations now converging; within a factor of 2 of the measurement result
 - Dispersion dependence needs study
 - Transverse coupling effect needs study
- **Measured bunch length exceeding those of Fokker-Planck solver**
 - Malitsky's re-bucketing simulation: satellite beam contribution
 - Estimation using measured rms bunch size underestimates IBS growth!
- **Strange Fokker-Planck behavior at injection**
 - Under investigation – work to be done

Transverse coupling



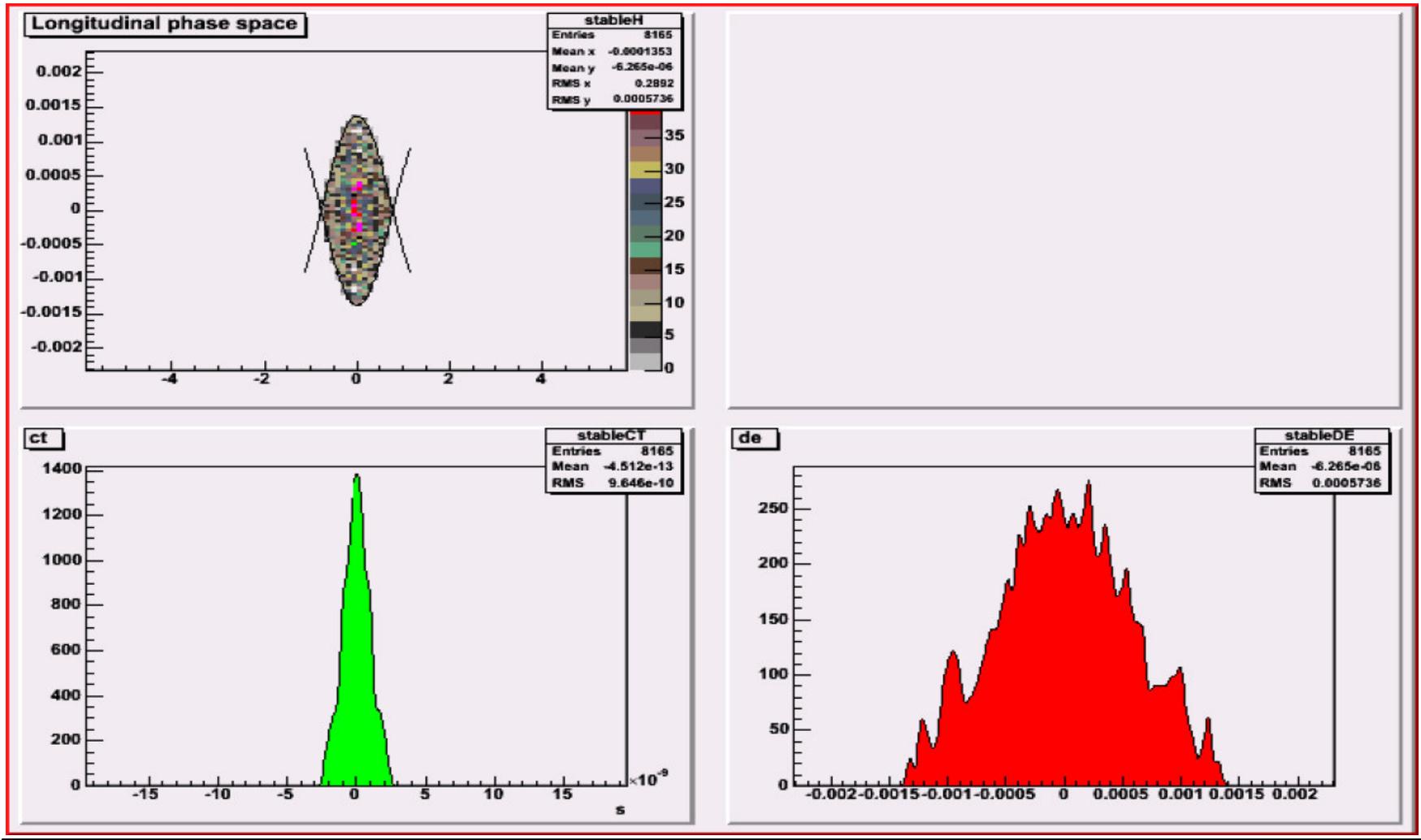
Re-bucketing with a full beam: satellites

- Large rms bunch length (2ns) due to satellite beams



Re-bucketing with a full beam

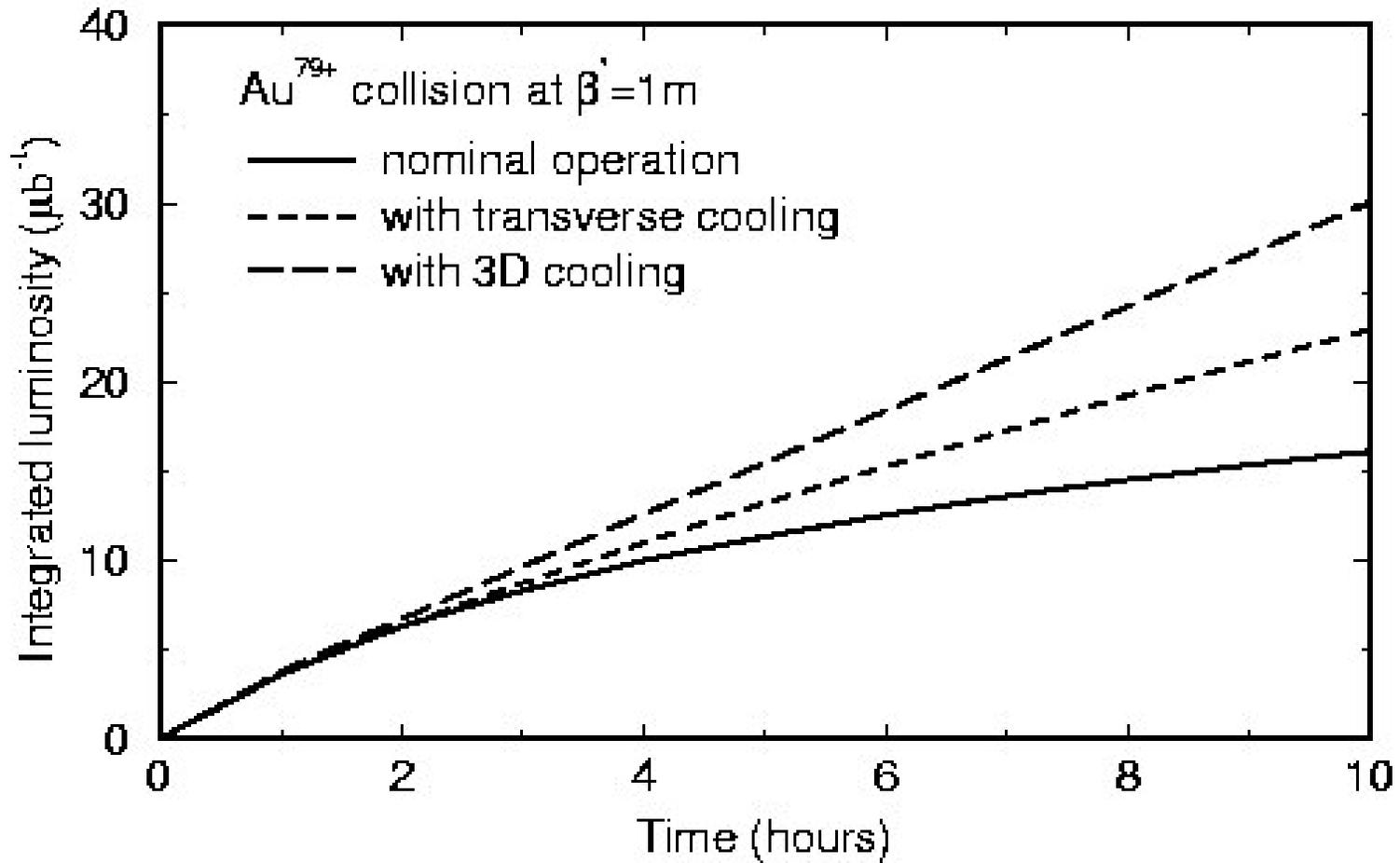
- Maximum rms bunch length is < 1 ns (bucket width $> 5 \sigma$)



Summary

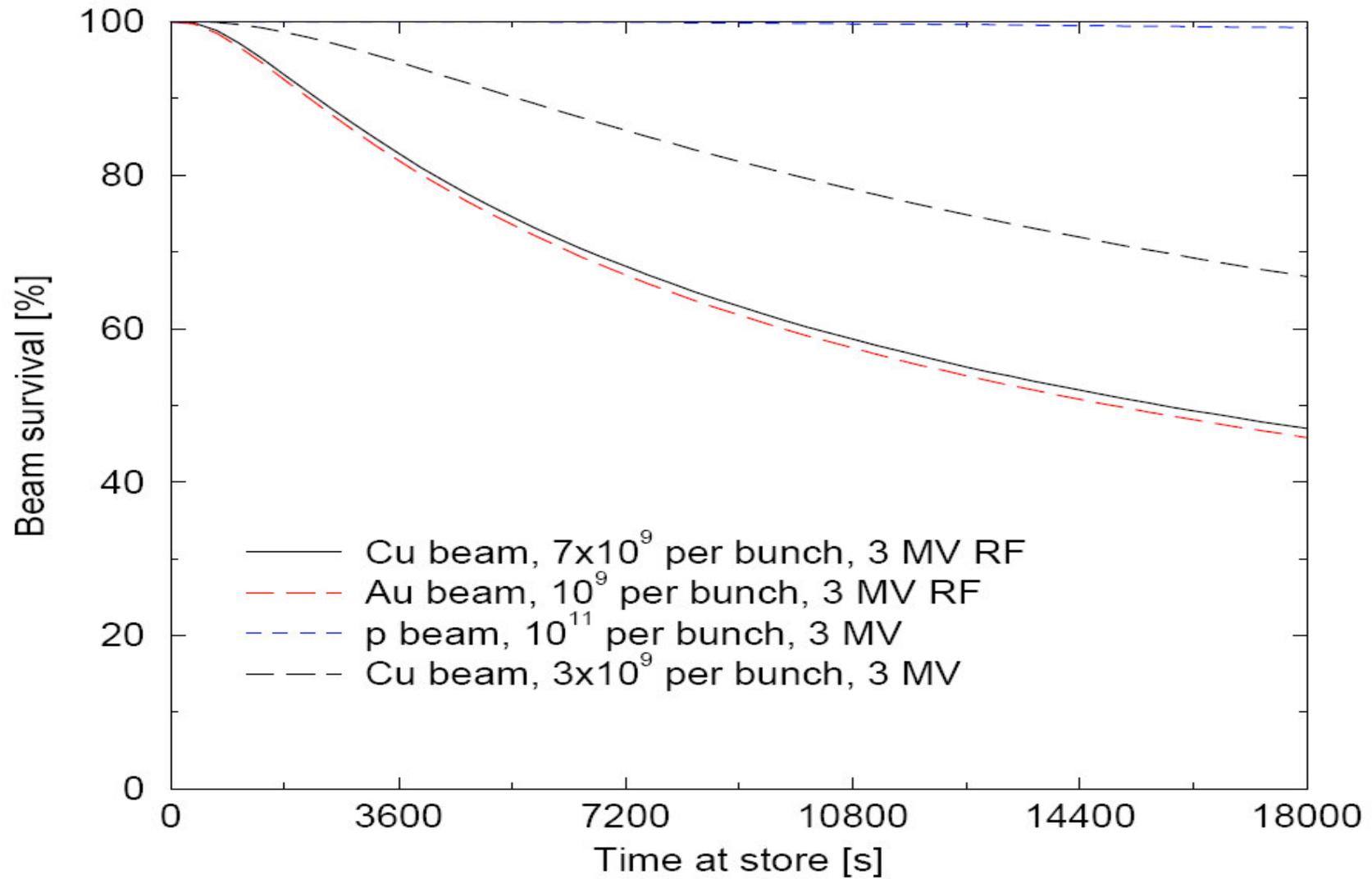
- The IBS benchmarking study at store is promising
 - Agreement on rms beam size growth is converging
 - » Longitudinal: agreement is within 10%
 - » Transverse: surely within a factor of 2
 - Agreement on de-bunching beam loss is good
 - Agreement on longitudinal profile is good
- Further studies needed
 - More accurate dispersion model & measurement
 - IBS under different transverse coupling conditions
 - IPM device calibration & lattice (beta-function) measurement
 - IBS study at injection

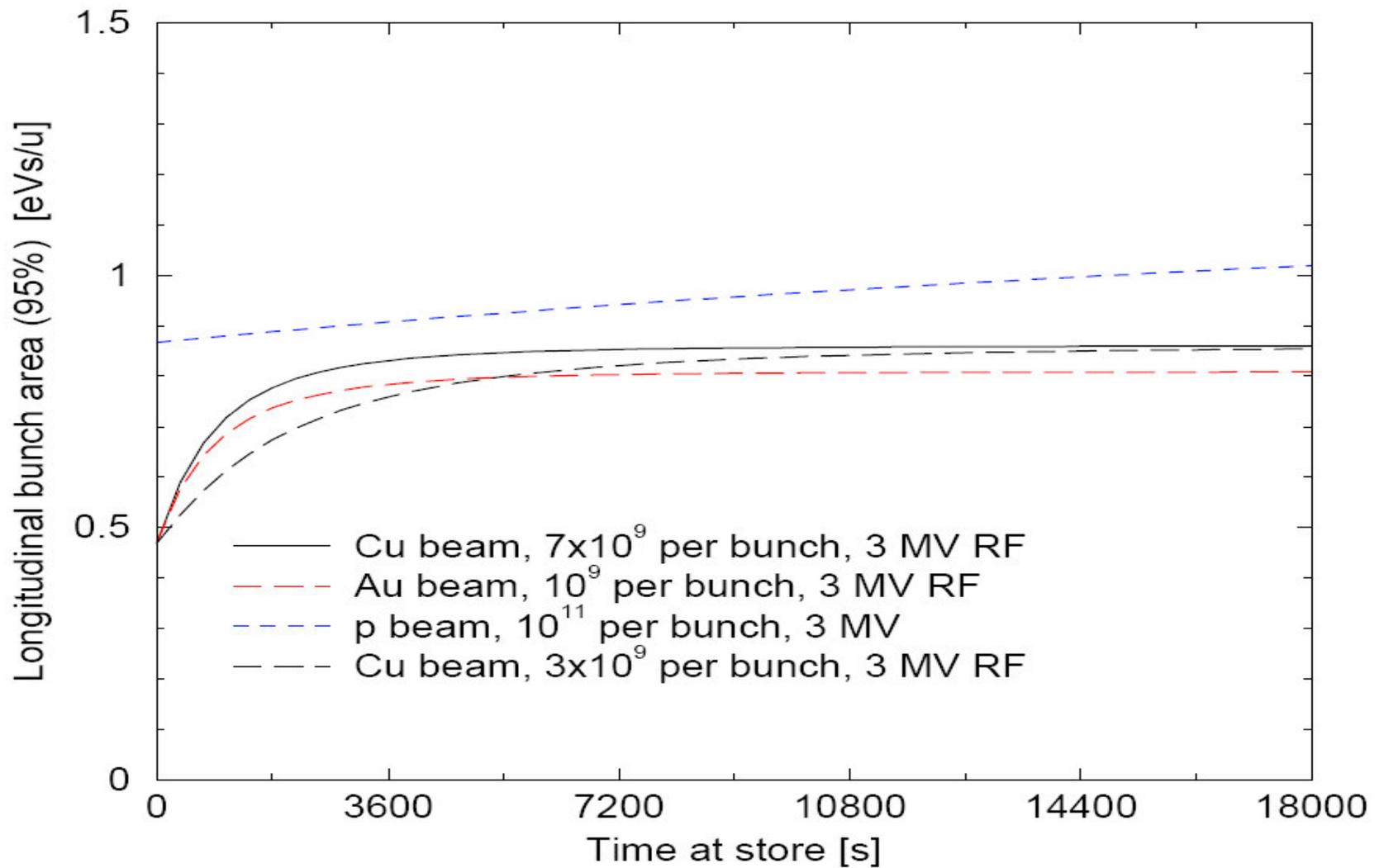
Predictions on beam cooling to counter IBS

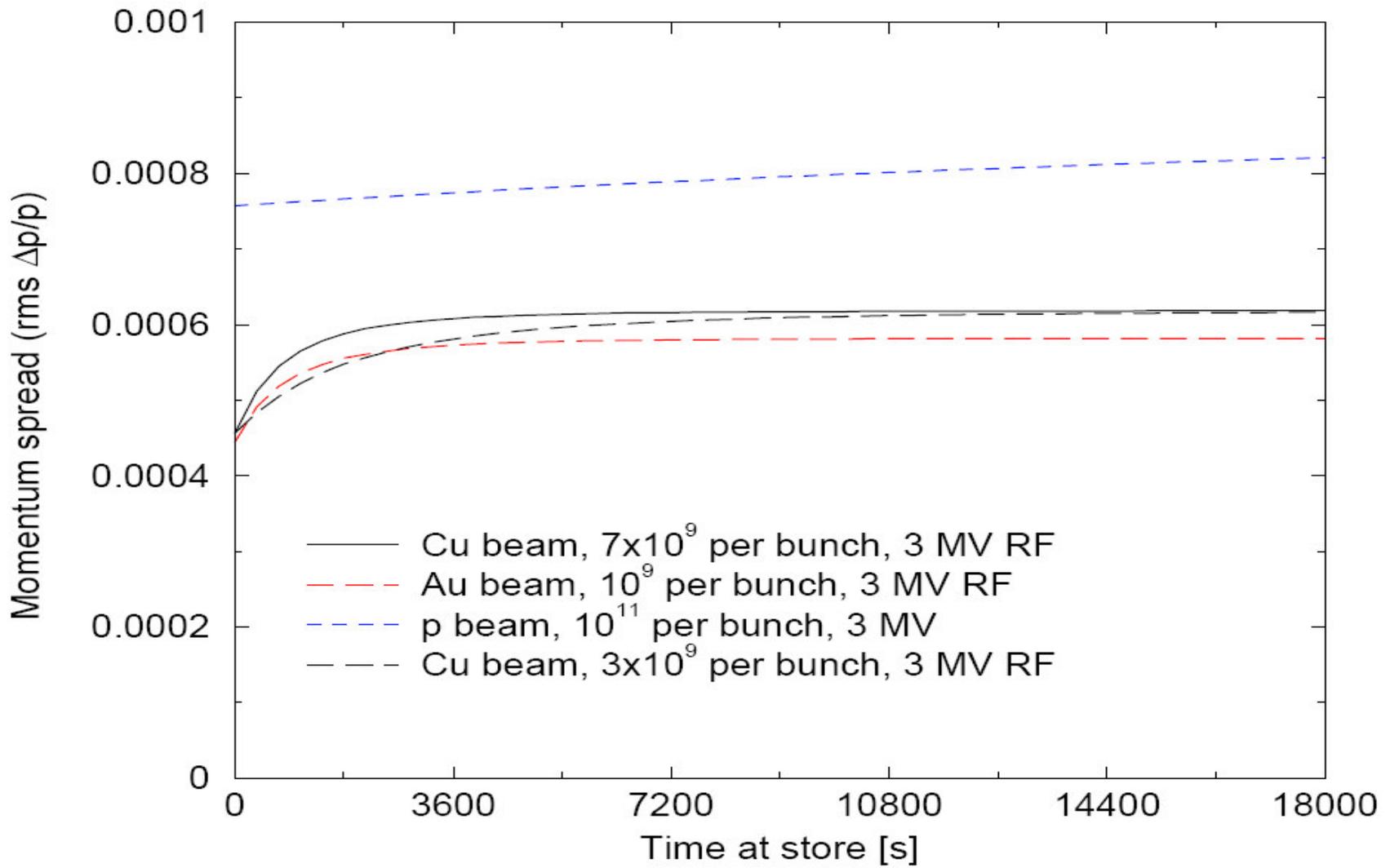


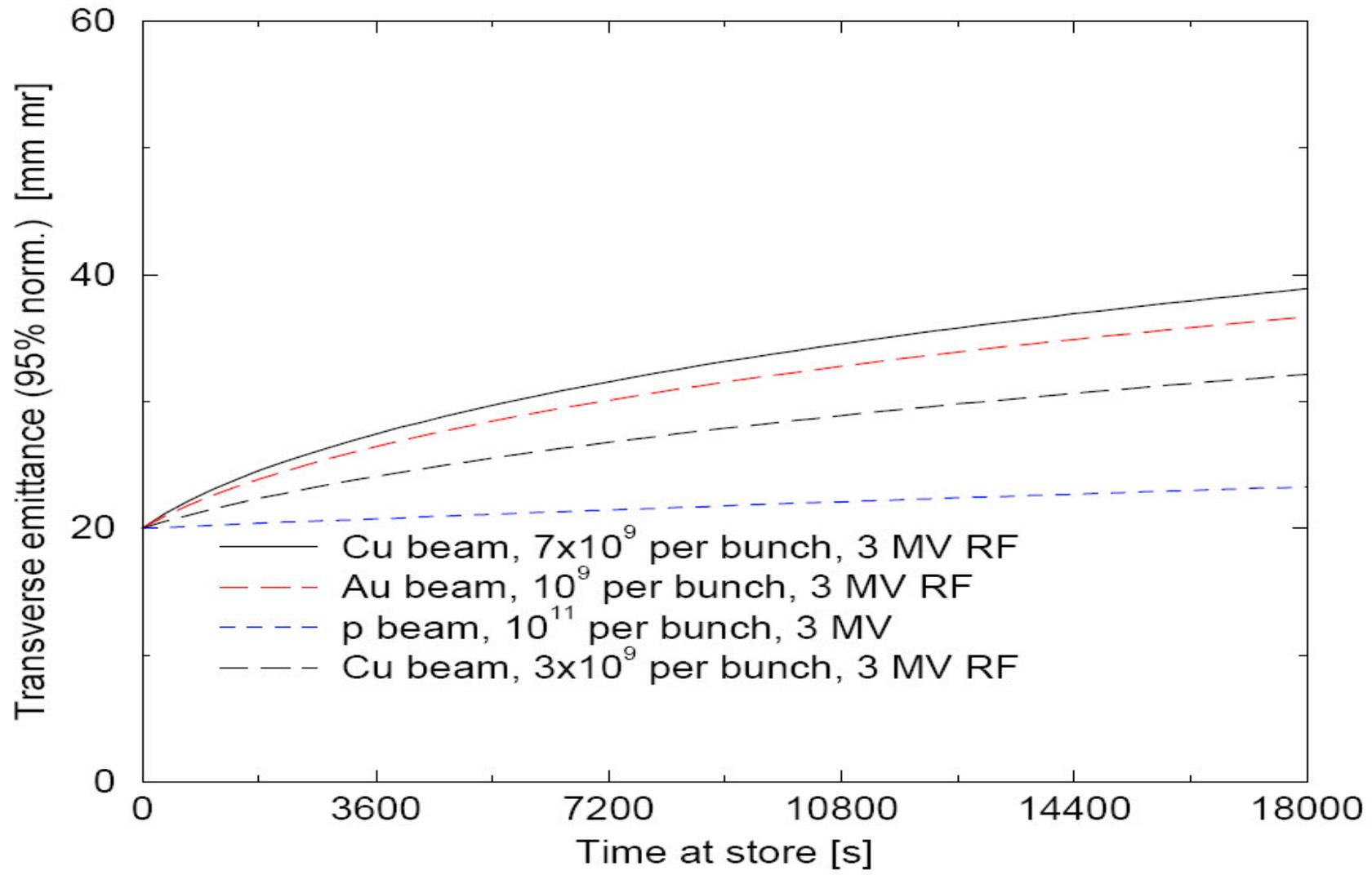
Predictions for the 2005 Cu run

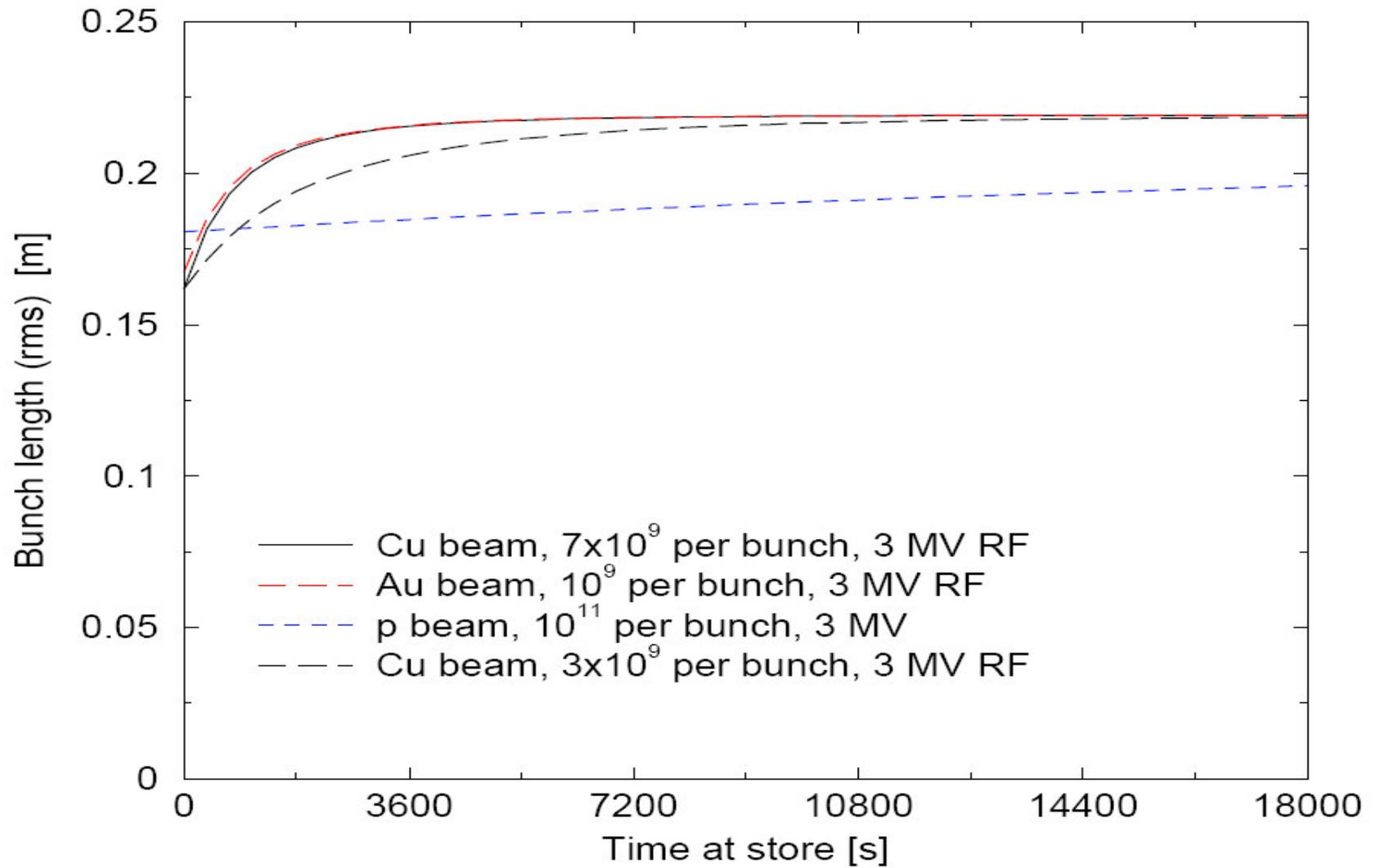
- Scale with NZ^4/A^2
- 6×10^9 Cu is equivalent to 1×10^9 Au
- IBS Expectation:
 - Similar behavior as Au store at 1×10^9 per bunch
 - Significant de-bunching expected; saturated bunch length and momentum spread
 - Usual rms growth prediction may not be adequate; BBFP can possibly better predict the IBS behavior











RHIC-5: estimates of IBS for Cu-Cu run

A. Fedotov and J. Wei (August 2004)

Parameters:

energy: 100 GeV

Cu: $Z=29$, $A=64$

95% normalized emittance: $20 \mu\text{m}$

$\beta^*=1\text{m}$

RF: $h=2520$, $V=3 \text{ MV}$

1. High-intensity case: $7 \cdot 10^9$ per bunch, 28 bunches
2. Low-intensity case: $3 \cdot 10^9$ per bunch, 45 bunches

Luminosity

- For high luminosities of Au ions, beam life time is limited by “burn-off” process of Au-Au collisions.
- Total cross section is about 212 barn (much higher than geometric cross section) due to
 - 1) Production of electron-positron pairs (117 barn)
 - 2) Coulomb dissociation of nucleus (95 barn)

For Cu: 1) is expected to be negligible 2) comparable to geometric cross section.

In simulation for Cu-Cu run only geometric cross section was assumed.

IBS Model

For IBS calculation of Cu-Cu run we used Jie Wei's model of IBS without high-energy approximation.

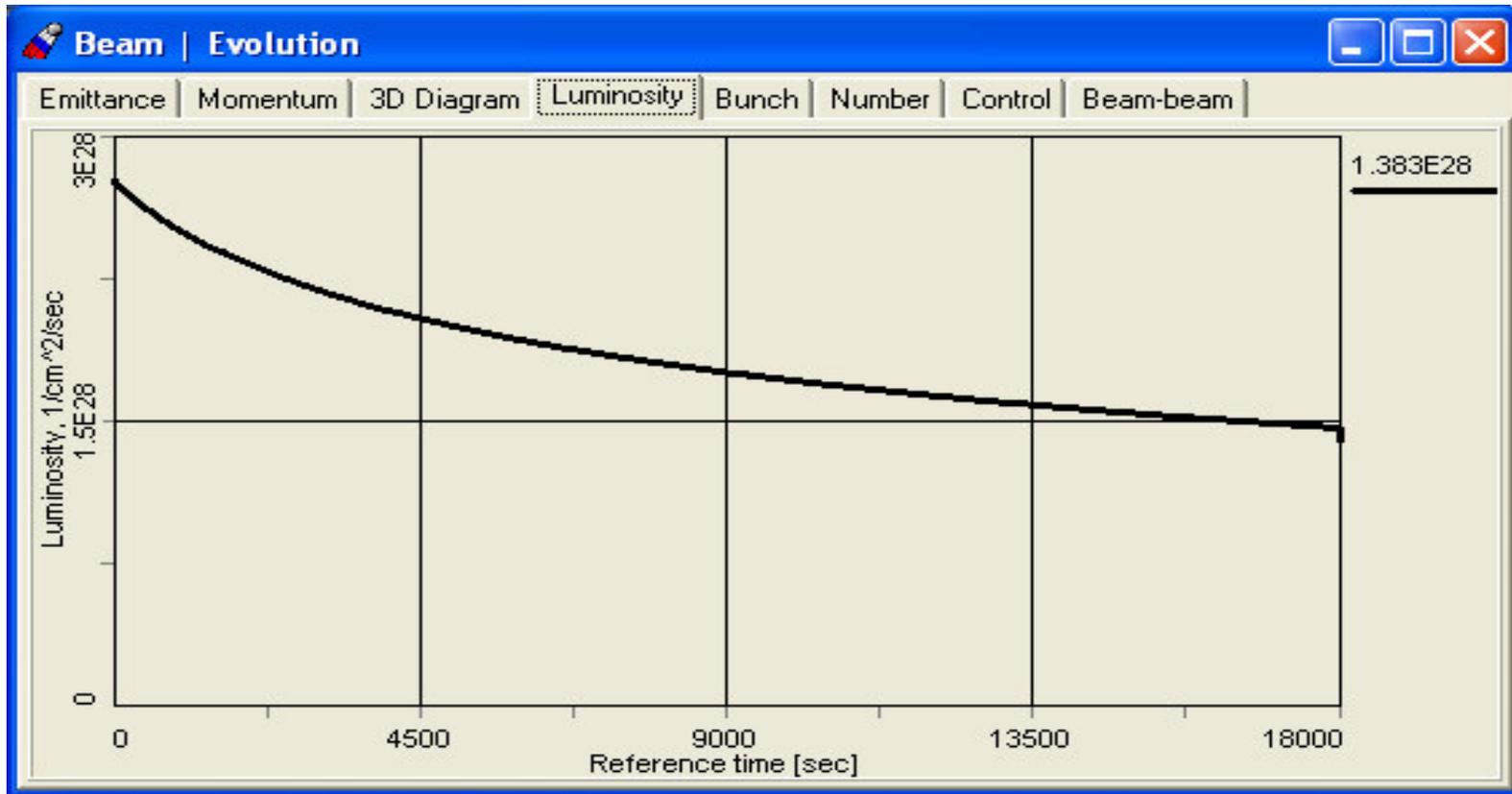
Simulations are done for:

1. Realistic RHIC lattice – under discussion.
2. FODO approximation – better describes emittance growth in RHIC (higher average dispersion) – used in present predictions for Cu run.

Simulations are done with the BETACOOOL and J. Wei's codes.

FODO approximation

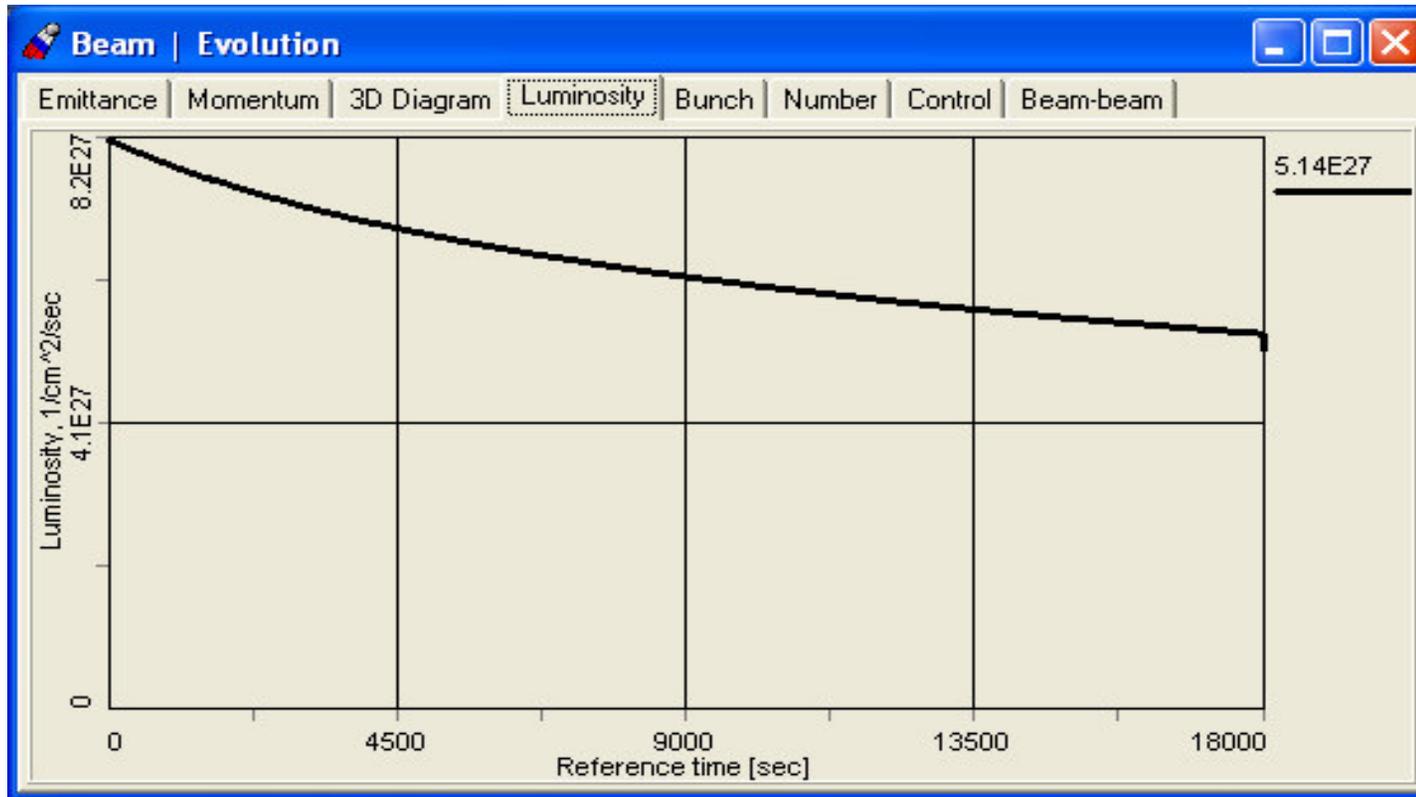
Cu-Cu, $N=7e9$, 28 bunches



Luminosity: $2.8e28 \rightarrow 1.4e28$ in 5 hours

FODO approximation

Cu-Cu, $N=3e9$, 45 bunches



Luminosity: $8.2e27 \rightarrow 5.1e27$ in 5 hours